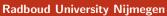
Process Theories and Graphical Language

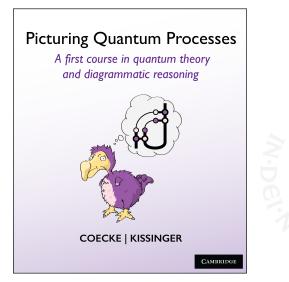
Aleks Kissinger

Institute for Computing and Information Sciences Radboud University Nijmegen

12th July 2016







Aleks Kissinger

Radboud University Nijmegen

Picturing Quantum Processes





Picturing Quantum Processes

When two systems [...] enter into temporary physical interaction due to known forces between them, [...] then they can no longer be described in the same way as before, viz. by endowing each of them with a representative of its own. I would not call that one but rather the characteristic trait of quantum mechanics, the one that enforces its entire departure from classical lines of thought.

- Erwin Schrödinger, 1935.



Picturing Quantum Processes

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- Erwin Schrödinger, 1935.

In quantum theory, *interaction* of systems is everything. **Diagrams** are the language of interaction.

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Picturing Quantum Processes

Q: How much of quantum theory can be understood just using diagrams and diagram transformation?

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Picturing Quantum Processes

Q: How much of quantum theory can be understood just using diagrams and diagram transformation?

A: Pretty much everything!



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Process theories and diagrams

Quantum processes

Classical and quantum interaction

Applications: a Hollywood-style trailer



Outline

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Process theories and diagrams

Quantum processes

Classical and quantum interaction

Applications: a Hollywood-style trailer







• A process is anything with zero or more *inputs* and zero or more *outputs*





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$$f(x,y) = x^2 + y$$





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• We could also write it like this:





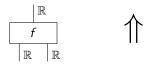
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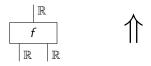
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• We could also write it like this:



• The labels on wires are called system-types or just types

More processes

• Similarly, computer programs are processes



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More processes

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- For example, a program that sorts lists might look like this:



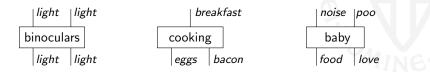
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More processes

- Similarly, computer programs are processes
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• These are also perfectly good processes:









• We can combine simple processes to make more complicted ones, described by diagrams:





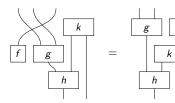


Diagrams

• We can combine simple processes to make more complicted ones, described by diagrams:



• The golden rule: only connectivity matters!





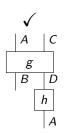
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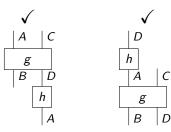
• Connections are only allowed where the types match, e.g.:







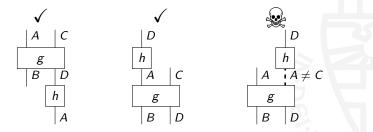
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Types

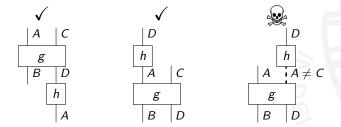
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Types

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• Types tell us when it makes sense to plug processes together

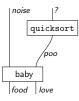
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Types and Process Theories

• Ill-typed diagrams are undefined:



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In fact, these processes don't ever make sense to plug together



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- A family of processes which <u>do</u> make sense together is called a process theory

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Types and Process Theories

• Ill-typed diagrams are undefined:



- In fact, these processes don't ever make sense to plug together
- A family of processes which <u>do</u> make sense together is called a process theory, e.g.
 - functions
 - linear maps
 - optical devices
 - proofs, ...

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Special processes: states and effects

• Processes with no inputs are called states:







Special processes: states and effects

Processes with no inputs are called states:

Interpret as: preparing a system in a particular configuration, where we don't care what came before.



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Special processes: states and effects

• Processes with no inputs are called states:

Interpret as: preparing a system in a particular configuration, where we don't care what came before.

• Processes with no outputs are called effects:



Interpret as: testing for a property π , where we don't care what happens after.





Numbers

• A number is a process with no inputs or outputs, written as:

 $\langle \lambda \rangle$ or just: λ





Numbers

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• Numbers always form a commutative monoid:

$$\langle \lambda \rangle \cdot \langle \mu \rangle := \langle \lambda \rangle \langle \mu \rangle$$
 1 :=

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Interpret as: what happens when a state meets an effect

effect
$$\left\{ \begin{array}{c} \swarrow \\ \downarrow \\ state \\ \psi \end{array} \right\}$$
 number

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 probability

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Interpret as: what happens when a state meets an effect, e.g.

effect
$$\left\{ \begin{array}{c} \swarrow \\ \bot \\ \downarrow \\ \end{array} \right\}$$
 probability state $\left\{ \begin{array}{c} \psi \\ \psi \end{array} \right\}$

This is called the (generalised) Born rule

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Process theories in general

Q: What kinds of behaviour can we study using just diagrams, and nothing else?

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Process theories in general

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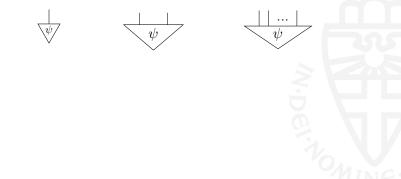
A: (Non-)separability

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Separable states

• States can be on a single system, two systems, or many systems:

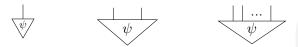


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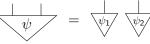


Separable states

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 A state ψ on two systems is ⊗-separable if there exist ψ₁, ψ₂ such that:

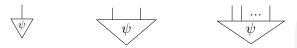


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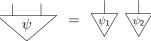


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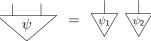


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 A state ψ on two systems is ⊗-separable if there exist ψ₁, ψ₂ such that:



- **Intuitively:** the properties of the system on the left are *independent* from those on the right
- In classical (deterministic) world, we expect all states to ⊗-separate

Characterising non-separability

...which is why non-separable states are way more interesting!





Characterising non-separability

- ...which is why non-separable states are way more interesting!
- But, how do we know we've found one?



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- i.e. that there do not exist states ψ_1, ψ_2 such that:

$$\psi$$
 = ψ



Characterising non-separability

- ...which is why non-separable states are way more interesting!
- But, how do we know we've found one?
- i.e. that there do not exist states ψ_1, ψ_2 such that:

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 = ψ_1 ψ_2

• Problem: Showing that something doesn't exist can be hard.



Characterising non-separability

Solution: Replace a negative property with a (stronger) postive one:



Characterising non-separability

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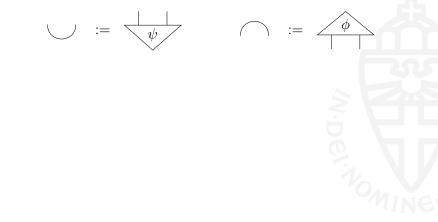
Definition

A state ψ is called *cup-state* if there exists an effect ϕ , called a *cap-effect*, such that:



Cup-states

• By introducing some clever notation:

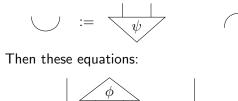


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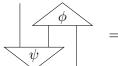
Process theories and diagrams Quantum processes Applications: a Hollywood-style trailer

Cup-states

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•



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:=

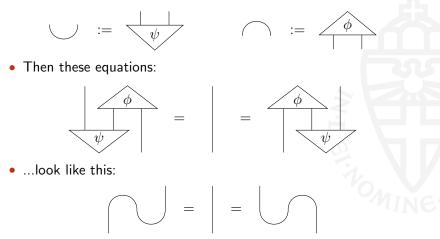
φ

 ψ

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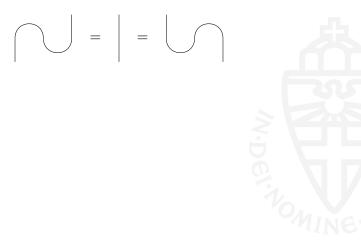
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Yank the wire!



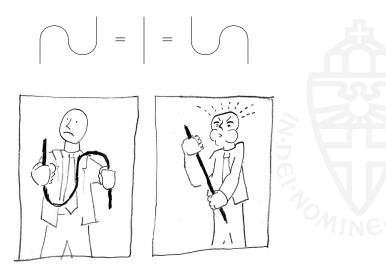
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Yank the wire!





A no-go theorem for separability

Theorem

If a process theory (i) has cup-states for every type and (ii) every state separates, then it is trivial.





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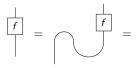
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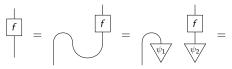
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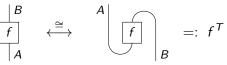
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Transpose



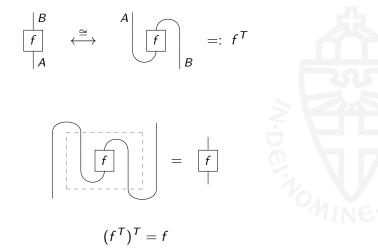




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Transpose



i.e.

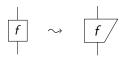
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Tranpose = rotation

A bit of a deformation:





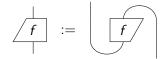
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f /

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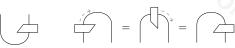
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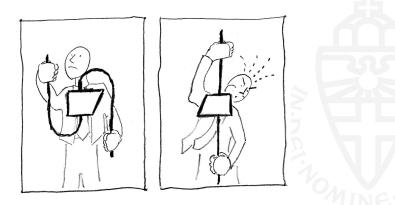




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Transpose = rotation



$\mathsf{Tranpose} = \mathsf{rotation}$

Specialised to states:

 $\widehat{\psi}$:= $\widehat{\psi}$

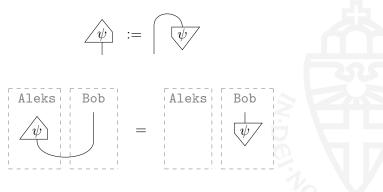


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Specialised to states:

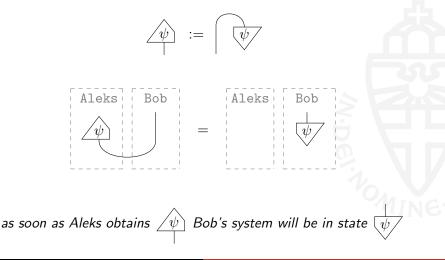


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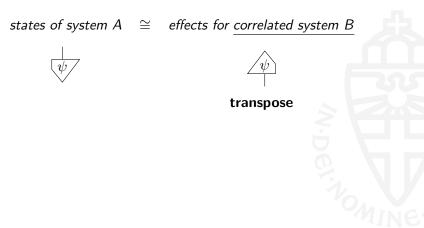
Specialised to states:



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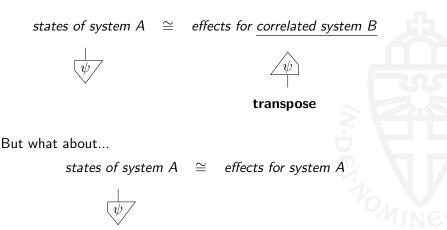
State/effect correspondence



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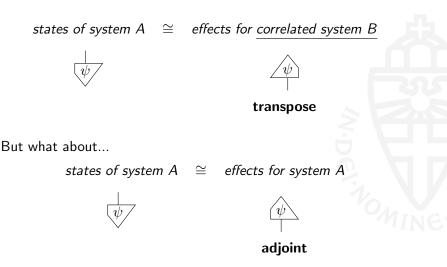
State/effect correspondence



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State/effect correspondence



Adjoints



state ψ

testing for ψ

U



Adjoints



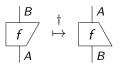


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state ψ

testing for ψ

Extends from states/effects to all processes:





Adjoints



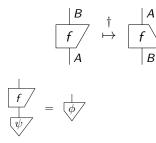


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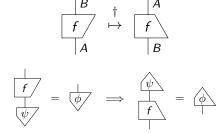




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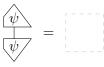






Normalised states and isometries

• Adjoints increase expressiveness, for instance can say when ψ is normalised:

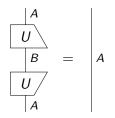




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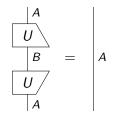




Normalised states and isometries

• Adjoints increase expressiveness, for instance can say when ψ is normalised:

• *U* is an *isometry*:



...and unitary, self-adjoint, positive, etc.

Conjugates

If we:



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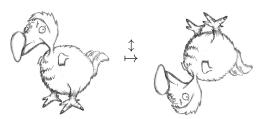


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Conjugates

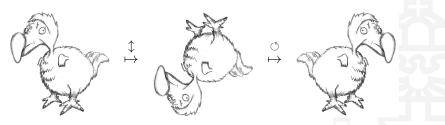
If we:





Conjugates

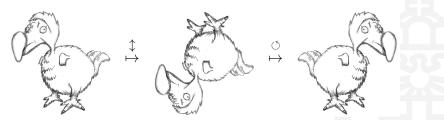
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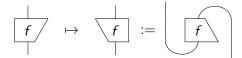
...we get horizontal reflection.

Conjugates

If we:



...we get horizontal reflection. The *conjugate*:



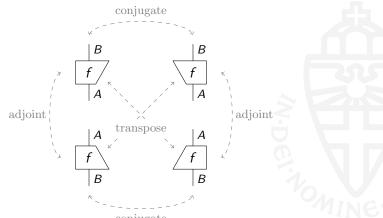
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4 kinds of box







conjugate

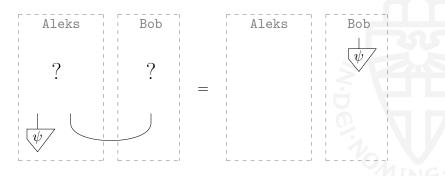
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Quantum teleportation: take 1

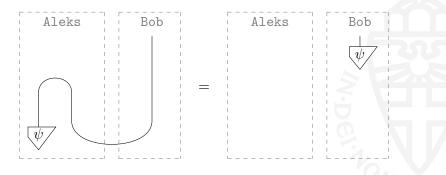
Can we fill in '?' to get this?





Quantum teleportation: take 1

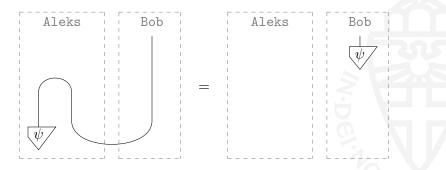
Here's a simple solution:





Quantum teleportation: take 1

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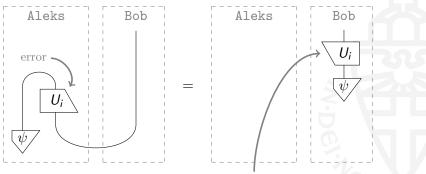


Problem: 'cap' can't be performed deterministically

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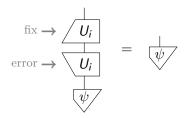
Bob's problem now!

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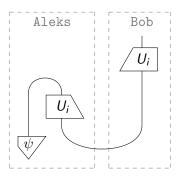


Quantum teleportation: take 1

Solution: Bob fixes the error.

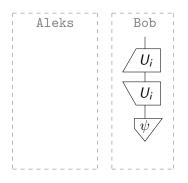


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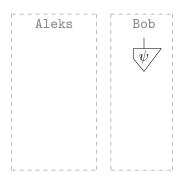




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Outline

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Process theories and diagrams

Quantum processes

Classical and quantum interaction

Applications: a Hollywood-style trailer



Hilbert space

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The starting point for quantum theory is the process theory of **linear maps**

Hilbert space

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The starting point for quantum theory is the process theory of **linear maps**, which has:

- **1** systems: Hilbert spaces
- Ø processes: complex linear maps

Hilbert space

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The starting point for quantum theory is the process theory of **linear maps**, which has:

- **1** systems: Hilbert spaces
- Ø processes: complex linear maps
- ...in particular, numbers are complex numbers.

Hilbert space

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Looking at the 'Born rule' for linear maps, we have a problem:

Hilbert space

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Looking at the 'Born rule' for linear maps, we have a problem:

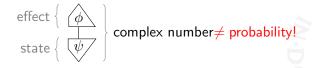


Hilbert space

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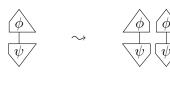


Looking at the 'Born rule' for linear maps, we have a problem:



Doubling

Solution: multiply by the conjugate:



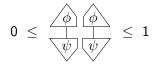


Doubling

Solution: multiply by the conjugate:



Then, for normalised ψ, ϕ :



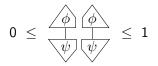


Doubling

Solution: multiply by the conjugate:

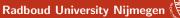


Then, for normalised ψ, ϕ :



(i.e. the 'usual' Born rule: $\overline{\langle \phi | \psi \rangle} \langle \phi | \psi \rangle = |\langle \phi | \psi \rangle|^2$)

Aleks Kissinger



Doubling

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New problem: We lost this:



Doubling

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New problem: We lost this:



...which was the basis of our interpretation for states, effects, and numbers.

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Doubling

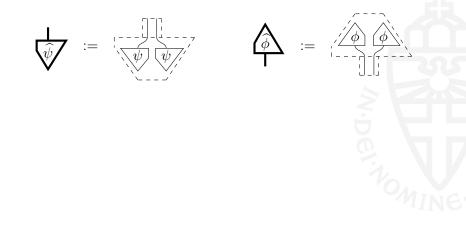
Solution: Make a new process theory with doubling 'baked in':

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Doubling

Solution: Make a new process theory with doubling 'baked in':

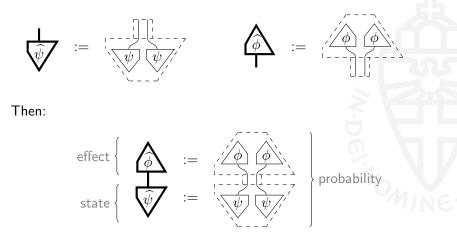


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Doubling

Solution: Make a new process theory with doubling 'baked in':







Doubling

The new process theory has doubled systems $\widehat{H} := H \otimes H$:

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Doubling

The new process theory has doubled systems $\widehat{H} := H \otimes H$:

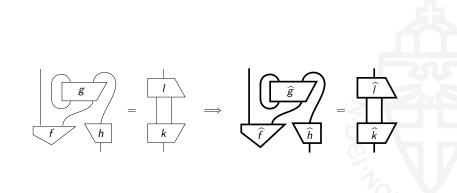
:= |||

and processes:

double $\begin{pmatrix} f \\ f \end{pmatrix}$:= $\begin{bmatrix} f \\ f \\ f \end{bmatrix}$ = $\begin{bmatrix} f \\ f \\ f \\ f \end{bmatrix}$

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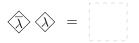
Doubling preserves diagrams



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...but kills global phases



(i.e. $\lambda = e^{i\alpha}$)

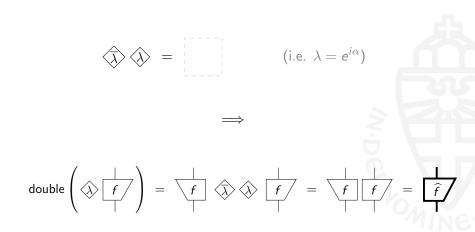


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...but kills global phases



Discarding

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Doubling also lets us do something we couldn't do before:





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Doubling also lets us do something we couldn't do before: throw stuff away!



Discarding

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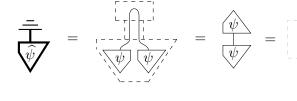


Doubling also lets us do something we couldn't do before: throw stuff away!

How? Like this:

Discarding

For normalised ψ , the two copies annihilate:

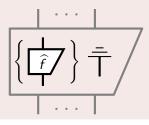




Quantum maps

Definition

The process theory of **quantum maps** has as types (doubled) Hilbert spaces \hat{H} and as processes:





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Two characterisations of 'pure'

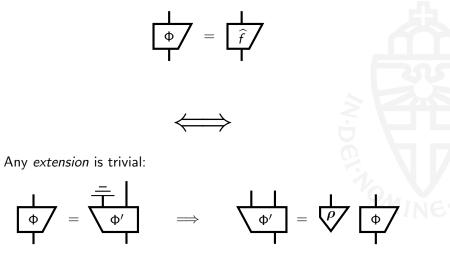
No discarding involved, i.e. for some f:

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Two characterisations of 'pure'

No discarding involved, i.e. for some f:



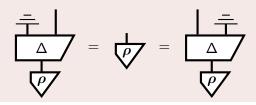
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Consequence: no-broadcasting

Theorem (No universal broadcasting)

There exists no quantum map Δ where:



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Consequence: no-broadcasting

Theorem (No universal broadcasting)

There exists no quantum map Δ where:

$$\begin{array}{c|c} - & - & - \\ \hline \end{array} \begin{array}{c} (1) \\ = \\ \end{array} \begin{array}{c} (1) \\ = \\ \end{array} \begin{array}{c} (1) \\ = \\ - \\ \hline \end{array} \begin{array}{c} - \\ - \\ - \\ \end{array} \end{array}$$
Proof. From (1):

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Consequence: no-broadcasting

Theorem (No universal broadcasting)

There exists no quantum map Δ where:

Proof. From (I):
$$\begin{array}{c} 1 \\ \Delta \\ 1 \end{array} = \begin{array}{c} 1 \\ P \\ P \\ P \end{array}$$

From (r):

$$= \begin{array}{c} 1 \\ \underline{-} \\ \underline{-} \\ 1 \end{array} = \begin{array}{c} \psi \\ \underline{-} \\ \underline{-} \\ \underline{-} \end{array}$$

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Consequence: no-broadcasting

Theorem (No universal broadcasting)

There exists no quantum map Δ where:

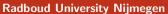
Proof. From (I):
$$\begin{array}{c} 1 \\ \Delta \\ 1 \end{array} = \begin{array}{c} 1 \\ P \\ P \\ P \end{array}$$

From (r):

$$=$$
 $\begin{bmatrix} -\frac{1}{2} \\ -\frac{1}{2} \end{bmatrix}$ $=$ $\begin{bmatrix} -\frac{1}{2} \\ -\frac{1}{2} \end{bmatrix}$

 \Rightarrow contradiction.







Aleks Kissinger

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Causality

A quantum map is called *causal* if:

$$\begin{bmatrix} - \\ T \\ 0 \\ T \end{bmatrix} = - T$$





Causality

A quantum map is called *causal* if:

$$\begin{bmatrix} \bar{\underline{-}} \\ \Phi \\ T \end{bmatrix} = \bar{\underline{-}}$$

If we discard the output of a process, it doesn't matter which process happened.





Causality

A quantum map is called *causal* if:

$$\begin{bmatrix} \bar{\underline{-}} \\ \Phi \\ T \end{bmatrix} = \bar{\underline{-}}$$

If we discard the output of a process, it doesn't matter which process happened.

 $\mathsf{causal} \iff \mathit{deterministically physically realisable}$



Consequence: no cap effect 🛞

Consequence: there is a unique causal effect, discarding:



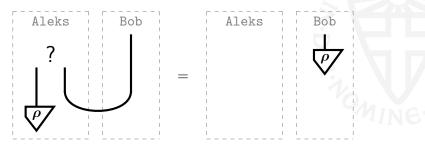
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Consequence: no cap effect 🛞

Consequence: there is a unique causal effect, discarding:



Hence 'deterministic quantum teleportation' must fail:



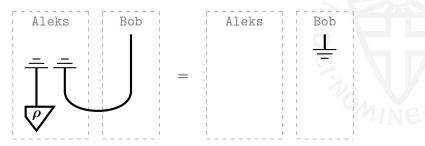


Consequence: no cap effect 🛞

Consequence: there is a unique causal effect, discarding:

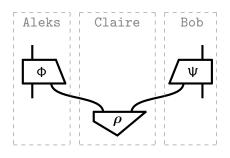


Hence 'deterministic quantum teleportation' must fail:





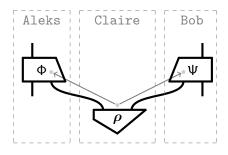
Process theories and diagrams Quantum processes Classical and quantum interaction Applications: a Hollywood-style trailer





Process theories and diagrams Quantum processes Classical and quantum interaction Applications: a Hollywood-style trailer

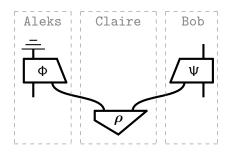
Consequence: no signalling 🙂



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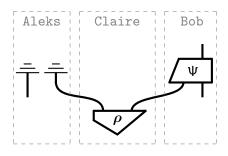


Process theories and diagrams Quantum processes Classical and quantum interaction Applications: a Hollywood-style trailer



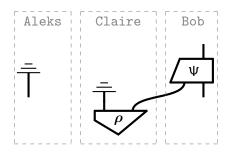


Process theories and diagrams Quantum processes Classical and quantum interaction Applications: a Hollywood-style trailer





Process theories and diagrams Quantum processes Classical and quantum interaction Applications: a Hollywood-style trailer



Outline

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Process theories and diagrams

Quantum processes

Classical and quantum interaction

Applications: a Hollywood-style trailer



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Double vs. single wires

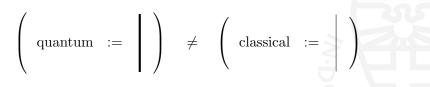
 $\left[\begin{array}{c} \text{quantum} := \end{array} \right]$



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Double vs. single wires



Classical values

$$\downarrow$$
 := 'providing classical value *i*'

Classical values

$$\stackrel{\perp}{i}$$
 := 'providing classical value *i*'

$$\frac{i}{1}$$
 := 'testing for classical value *i*'



Classical values

$$i$$
 := 'providing classical value *i*'

$$\stackrel{i}{\downarrow}$$
 := 'testing for classical value *i*'

$$\begin{array}{c} \underbrace{j}\\ \hline \\ i \end{array} = \begin{cases} 1 & \text{if } i = j \\ 0 & \text{if } i \neq j \end{cases}$$

Aleks Kissinger

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Classical values

$$i$$
 := 'providing classical value *i*'

$$\stackrel{\frown}{i}$$
 := 'testing for classical value *i*'

$$\begin{array}{c} \overbrace{j}\\ \hline \\ \hline \\ i \end{array} = \begin{cases} 1 & \text{if } i = j \\ 0 & \text{if } i \neq j \end{cases}$$

$$(\Rightarrow \text{ONB})$$



Classical states

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General state of a classical system:

$$\bigvee_{p} := \sum_{i} p_{i} \bigvee_{i}$$

probability distributions

Classical states

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General state of a classical system:

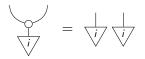
 \bigvee_{i}^{p} := $\sum_{i} p_{i} \bigvee_{i}^{l} \leftarrow$ probability distributions

Hence:

$$\bigvee_{i}^{\perp}$$
 \leftarrow point distributions

Copy and delete

Unlike quantum states, classical values can be copied:

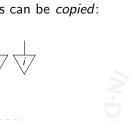




Copy and delete

Unlike quantum states, classical values can be copied:

and *deleted*:

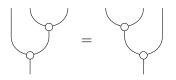




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Copy and delete

These satisfy some equations you would expect:



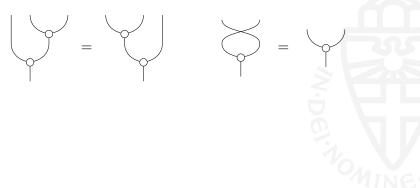


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These satisfy some equations you would expect:

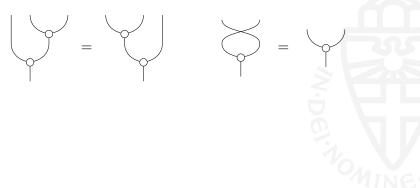


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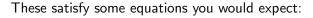
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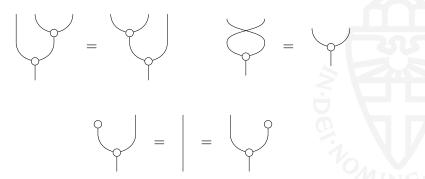


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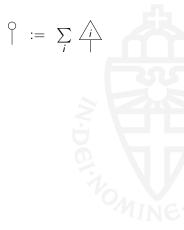


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Other classical maps

 $\overline{\vee}$ $:= \sum_{i} \bigvee_{i}^{i}$

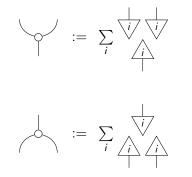


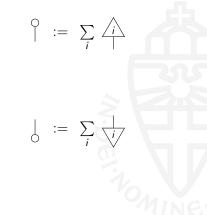
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Other classical maps

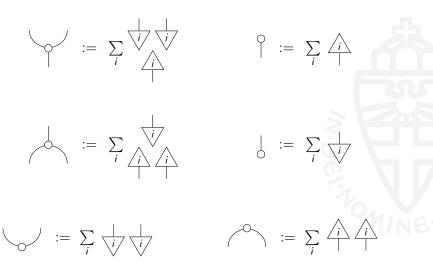




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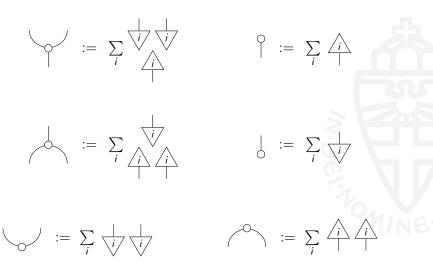
Other classical maps



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Other classical maps



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...satisfying lots of equations

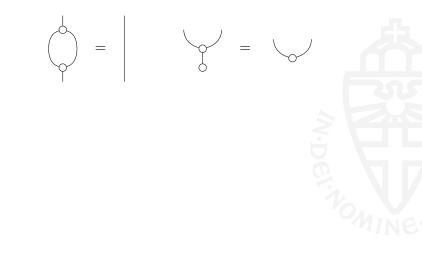




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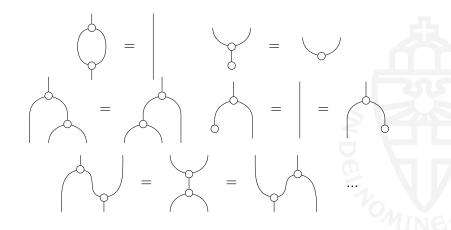
...satisfying lots of equations



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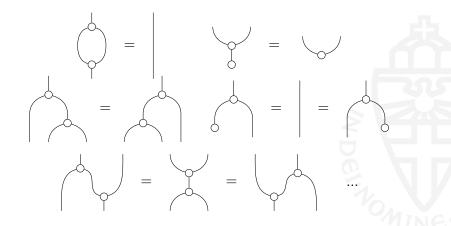
...satisfying lots of equations



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...satisfying lots of equations



When does it end???

Spiders



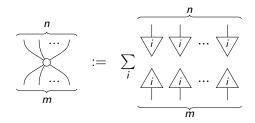


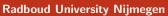
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Spiders

All of these are special cases of *spiders*:

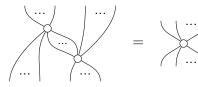






Spiders

The only equation you need to remember is this one:

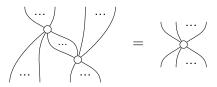






Spiders

The only equation you need to remember is this one:

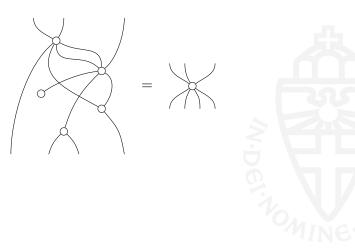


When spiders meet, they fuse together.

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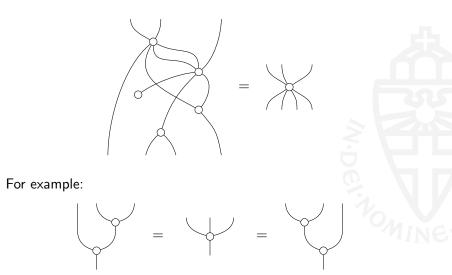
Spider reasoning



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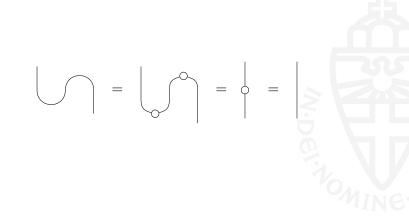
Spider reasoning



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Spider reasoning \Rightarrow string diagram reasoning





How do we recognise spiders?

Suppose we have something that 'behaves like' a spider:





How do we recognise spiders?

Suppose we have something that 'behaves like' a spider:



Do we know it is one?

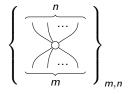
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Spiders = 'diagrammatic ONBs'

Yes!



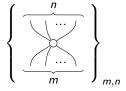


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Spiders = 'diagrammatic ONBs'

Yes!





i

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Classical and quantum interaction





Classical and quantum interaction

Classical values can be encoded as quantum states, via doubling:



Classical and quantum interaction

Classical values can be encoded as quantum states, via doubling:

This is our first classical-quantum map, encode.



Classical and quantum interaction

Classical values can be encoded as quantum states, via doubling:

This is our first classical-quantum map, *encode*. It's a copy-spider in disguise:

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Measuring quantum states

The adjoint of *encode* is *measure*:



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Measuring quantum states

The adjoint of *encode* is *measure*:

quantum state $\langle \bigvee^{\text{probability distribution}} \rangle$

This represents measuring w.r.t.

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Measuring quantum states

The adjoint of *encode* is *measure*:

quantum state $\left\{ \begin{array}{c} & & \\$

This represents measuring w.r.t.

...where probabilities come from the Born rule:

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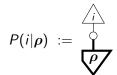
Measuring quantum states

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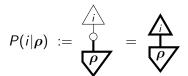
Measuring quantum states

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quantum state $\left\{ \begin{array}{c} & & \\$

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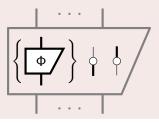
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Classical-quantum maps

Definition

The process theory of **cq-maps** has as processes diagrams of quantum maps and encode/decode:



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Quantum processes

Causality generalises to cq-maps:

$$\begin{bmatrix} 0 & -\frac{1}{2} \\ 0 & -\frac{1}{2} \\ 0 & -\frac{1}{2} \end{bmatrix} = 0 \begin{bmatrix} -\frac{1}{2} \\ -\frac{1}{2} \end{bmatrix}$$

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Quantum processes

Causality generalises to cq-maps:

$$\begin{bmatrix} \phi & -\frac{1}{2} \\ \phi \\ \phi \end{bmatrix} = \begin{pmatrix} \phi & -\frac{1}{2} \\ 0 &$$

quantum processes := causal cq-maps



Special case: quantum measurements

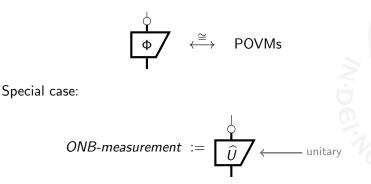
A *measurement* is any **quantum process** from a quantum system to a classical one:

$$\begin{array}{c} & & \\ & & \\ \hline \Phi \end{array} & \stackrel{\cong}{\longleftrightarrow} & \mathsf{POVMs} \end{array}$$



Special case: quantum measurements

A *measurement* is any **quantum process** from a quantum system to a classical one:



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Special case: controlled-operations

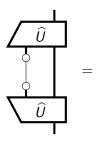
A **quantum process** with a classical input is a *controlled operation*:





Special case: controlled-operations

A controlled isometry furthermore satisfies:







Special case: controlled-operations

Suppose we can use a single \hat{U} to build a *controlled isometry*:



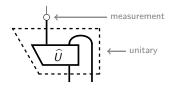


Special case: controlled-operations

Suppose we can use a single \hat{U} to build a *controlled isometry*:



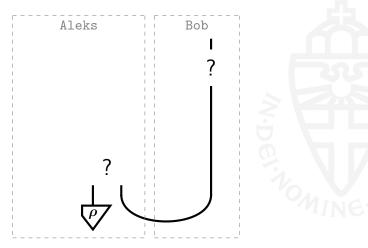
...and an ONB measurement:



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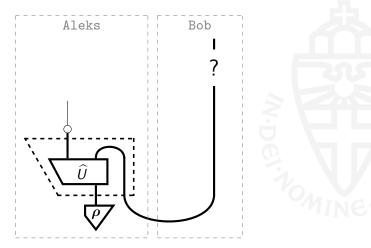
Quantum teleportation: take 2



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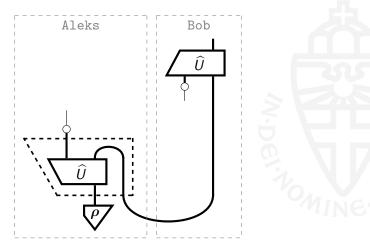
Quantum teleportation: take 2



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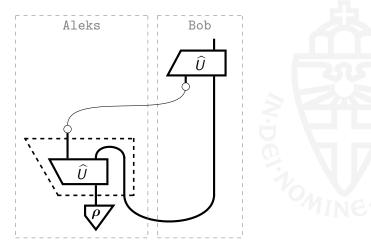
Quantum teleportation: take 2



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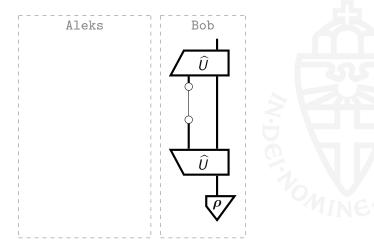
Quantum teleportation: take 2



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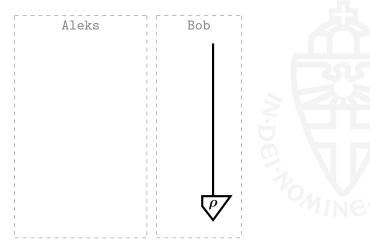
Quantum teleportation: take 2



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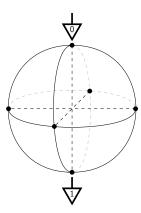
Quantum teleportation: take 2



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Complementary bases

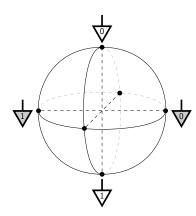




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Complementary bases



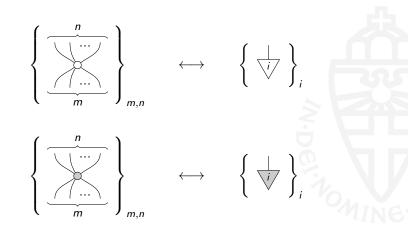
Aleks Kissinger

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Complementary bases



Complementarity



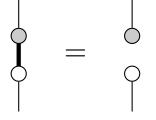
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Complementarity



Interpretation:

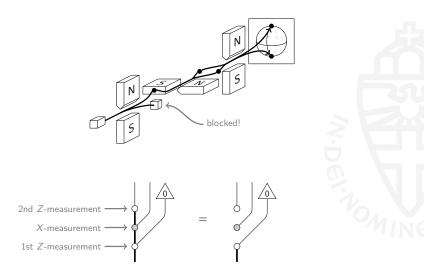
(encode in \bigcirc) THEN (measure in \bigcirc) = (no data flow)



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Consequence: Stern-Gerlach



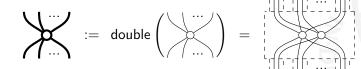
Process Theories and Graphical Language

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Quantum computation

Doubling a classical spider gives a *quantum spider*:

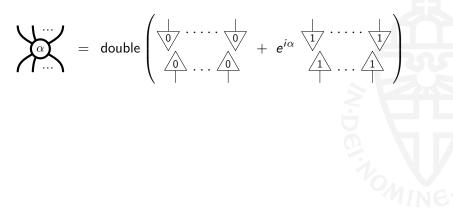


Universality

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By decorating quantum spiders with phases:

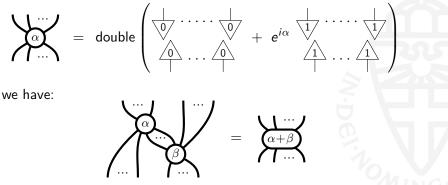


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Universality

By decorating quantum spiders with phases:



and spider-diagrams become universal for quantum computation!



Soundness and completeness

Restricting the phase group to $\mathbb{Z}_4 \cong \{0, \frac{\pi}{2}, \pi, \frac{3\pi}{2}\} \subset U(1)$ gives *stabiliser QT*.

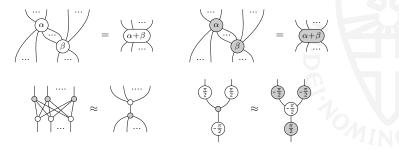
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Soundness and completeness

Restricting the phase group to $\mathbb{Z}_4 \cong \{0, \frac{\pi}{2}, \pi, \frac{3\pi}{2}\} \subset U(1)$ gives *stabiliser QT*.

Sound and complete presentation via the ZX-calculus:



Outline

Process theories and diagrams

Quantum processes

Classical and quantum interaction

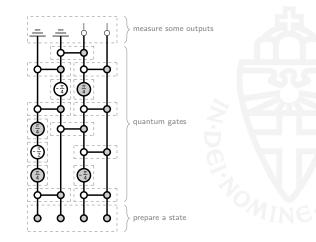
Applications: a Hollywood-style trailer





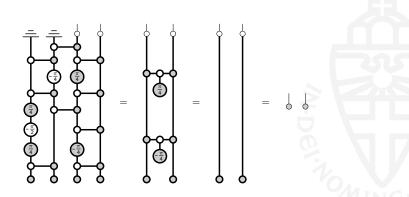
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Quantum circuits and rewriting



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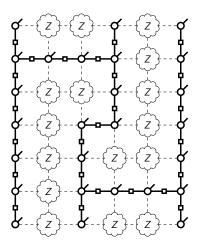
Quantum circuits and rewriting



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Measurement-based quantum computing

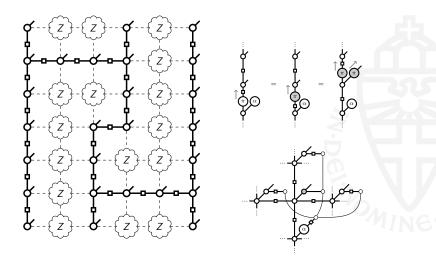




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Measurement-based quantum computing

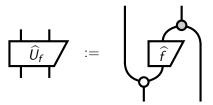


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Quantum algorithms

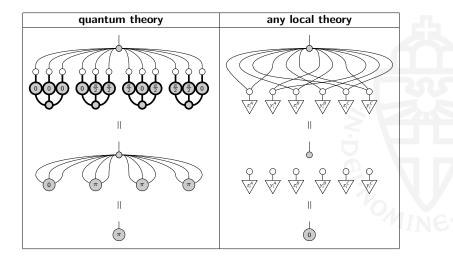
Spiders can be used to build quantum oracles:



 \Rightarrow simple derivations of **Deutsch-Jozsa**, **quantum seach**, and **hidden subgroup** algorithms.

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GHZ/Mermin non-locality

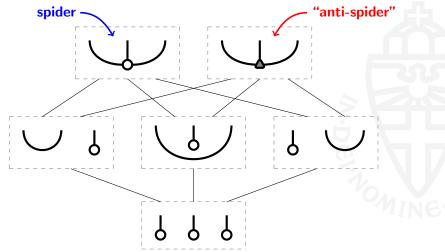


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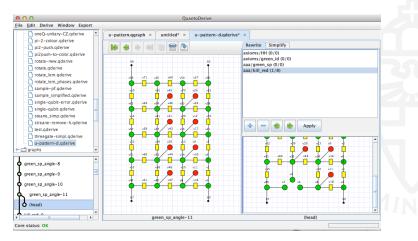
Multipartite entanglement

SLOCC-classification of 3 qubits:



Automation

Quantomatic:



Aleks Kissinger

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Thanks! Joint work with Bob Coecke (book):



...and many more!



Abramsky, Backens, Duncan, Edwards, Gogioso, Hadzihasanovic, Heunen, Lal, Merry, Pavlovic, Perdrix, Quick, Selinger, Vicary, Zamdzhiev, ...

http://quantomatic.github.io

Aleks Kissinger